

AOS 630: Introduction to Atmospheric  
and Oceanic Physics  
Lecture 3 Fall 2021  
*Moist air and hydrostatic balance*

Ángel F. Adames-Corraliza  
angel.adamescorraliza@wisc.edu



# Announcements

Homework 1 is now online. It is due two weeks from today.

# Last Class: the equation of state

For dry air

$$p_d = \rho_d R_d T$$

$$R_d = 287 \mathbf{J\ kg^{-1}\ K^{-1}}$$

Dry gas constant

Partial pressure for water vapor

$$e = \rho_v R_v T$$

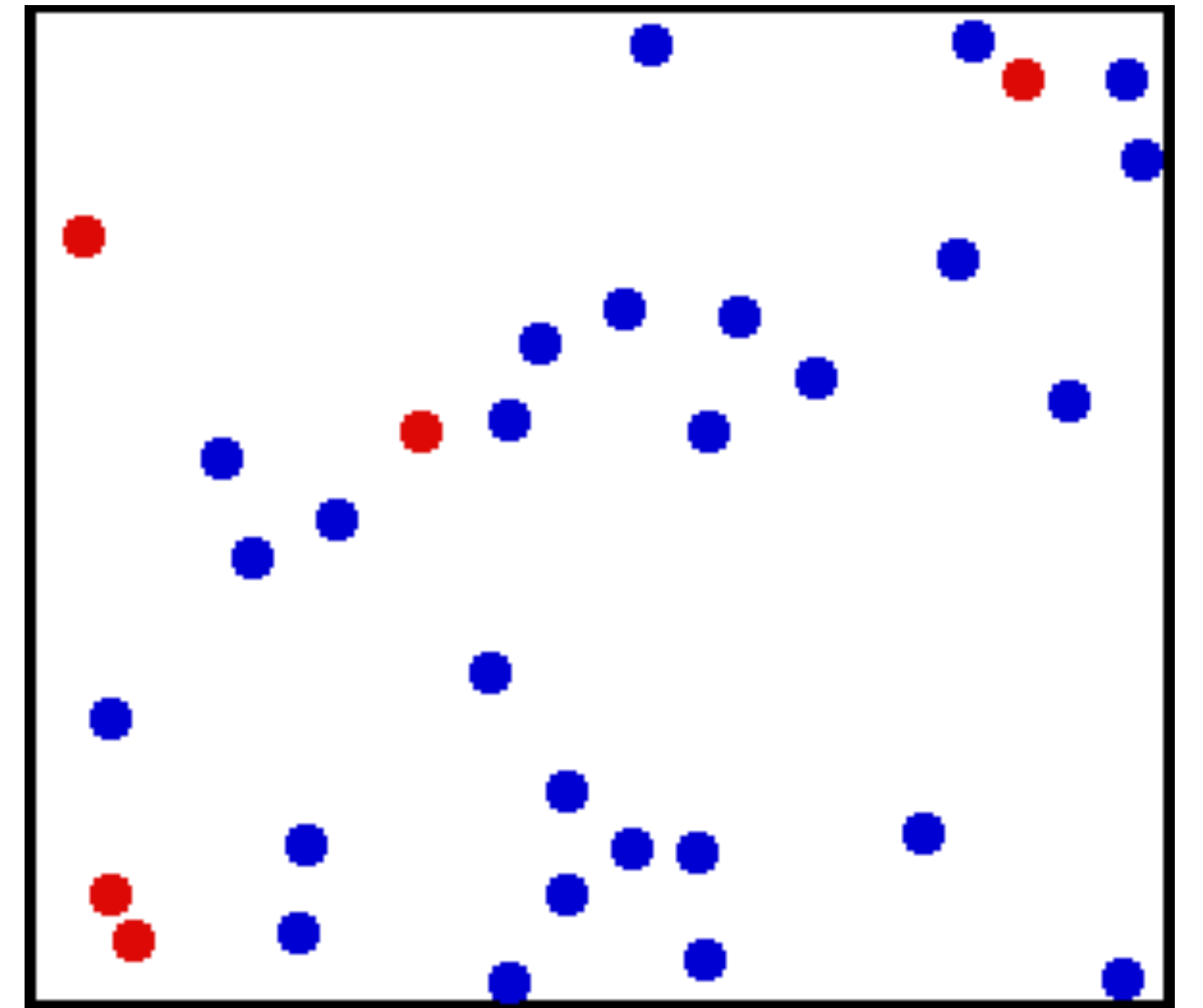
$$R_v = 461 \mathbf{J\ kg^{-1}\ K^{-1}}$$

Water vapor constant

Partial pressure for moist air

$$p = p_d + e$$

$$p = (\rho_d R_d + \rho_v R_v) T$$





# Equation of state for seawater

$$\rho = \rho_0 \left[ 1 - \beta_T (T - T_0) + \beta_s (S - S_0) + \beta_p (p - p_0) \right]$$

Symbol	Description	Value
$\rho_0$	reference density	$1.027 \times 10^3 \text{ kg m}^{-3}$
$\alpha_0$	reference specific volume	$9.738 \times 10^{-4} \text{ m}^3 \text{ kg}^{-1}$
$T_0$	reference temperature	283 K
$S_0$	reference salinity	35 ppt = $35 \text{ g kg}^{-1}$
$c_{s0}$	reference sound speed	$1490 \text{ m s}^{-1}$
$\beta_T$	thermal expansion coefficient	$1.67 \times 10^{-4} \text{ K}^{-1}$
$\beta_s$	haline contraction coefficient	$0.78 \times 10^{-3} \text{ ppt}^{-1}$
$\beta_p$	compressibility coefficient ( $= \alpha_0/c_{s0}^2$ )	$4.39 \times 10^{-10} \text{ m s}^2 \text{ kg}^{-1}$
$c_{p0}$	specific heat capacity at const. pressure	$3986 \text{ J kg}^{-1} \text{ K}^{-1}$

This is just an approximation.  
The equation that is used in  
ocean modeling is much more  
complicated!





# Today

Examine how water vapor modifies  
the equation of state

Introduce the hydrostatic equation



# Supplementary reading

Petty

Sections 3.4 and Chapter 4

Wallace and Hobbs

Section 3.1.1 and 3.2



# Water vapor

Water vapor is roughly an ideal gas. It follows Dalton's law of partial pressures (the total pressure is the sum of the pressure of all the constituent gases).

$$e\alpha_v = R_v T$$

The mixing ratio is the amount of water vapor mass per unit of dry air

$$r_v = \frac{m_v}{m_d}$$

The specific humidity is the amount of water vapor per unit of total air mass.

$$q_v = \frac{m_v}{m_d + m_v} \quad q_v \simeq r_v$$



# Water vapor

Using the ideal gas law we can express the mixing ratio and specific humidity in terms of pressure

$$e = \rho_v R_v T \qquad p = \rho R_d T$$

Which are written as

$$r_v \simeq \varepsilon \frac{e}{p} \qquad q_v \simeq \varepsilon \frac{e}{p + e}$$

$$\varepsilon = R_d / R_v = 0.622$$



# Modifications of the equations due to moisture

Water vapor is lighter than the dry air molecules  $O_2$  and  $N_2$

This means that a volume of humid air that is at the same temperature as a volume of dry air is actually less dense

<b>Constituent<sup>a</sup></b>	<b>Molecular weight</b>	<b>Fractional concentration by volume</b>
Nitrogen ( $N_2$ )	28.013	78.08%
Oxygen ( $O_2$ )	32.000	20.95%
Argon (Ar)	39.95	0.93%
<b>Water vapor (<math>H_2O</math>)</b>	18.02	0–5%
<b>Carbon dioxide (<math>CO_2</math>)</b>	44.01	380 ppm
Neon (Ne)	20.18	18 ppm
Helium (He)	4.00	5 ppm
<b>Methane (<math>CH_4</math>)</b>	16.04	1.75 ppm
Krypton (Kr)	83.80	1 ppm
Hydrogen ( $H_2$ )	2.02	0.5 ppm
<b>Nitrous oxide (<math>N_2O</math>)</b>	56.03	0.3 ppm
<b>Ozone (<math>O_3</math>)</b>	48.00	0–0.1 ppm



To take into account this change in density we define the virtual temperature

$$p\alpha = R_d T_v$$

$$T_v \simeq T (1 + 0.61 q_v)$$

The virtual temperature is the temperature dry air would have if it had the same density as the moist air at the same pressure.



# Exercise

Calculate the virtual temperature for the locations below.

Based on your answer, do you think the virtual temperature correction may be important somewhere and why?

$$T_v \simeq T (1 + 0.61q_v)$$

Location	Temperature (°C)	Mixing ratio (g/kg)
Utqiagvik (Barrow), AK	0	5
Gaylord, MI	15	10
Singapore	30	30

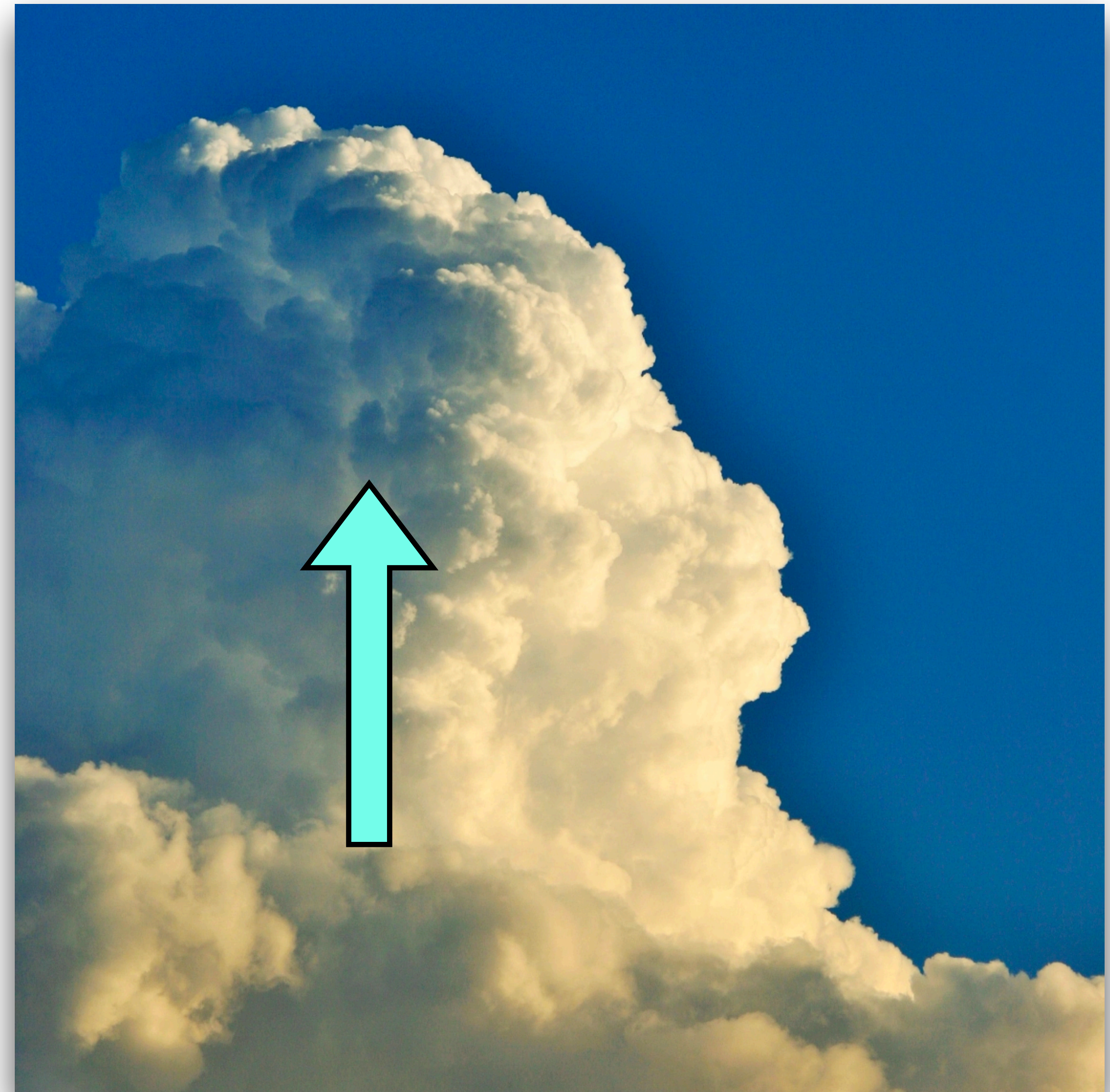


Newton's second law dictates that acceleration must result from a net sum of forces.

Apply this to vertical motion

$$\frac{Dw}{Dt} = \frac{1}{m} \sum_i F_z$$

$$\frac{D}{Dt} = \frac{d}{dt}$$





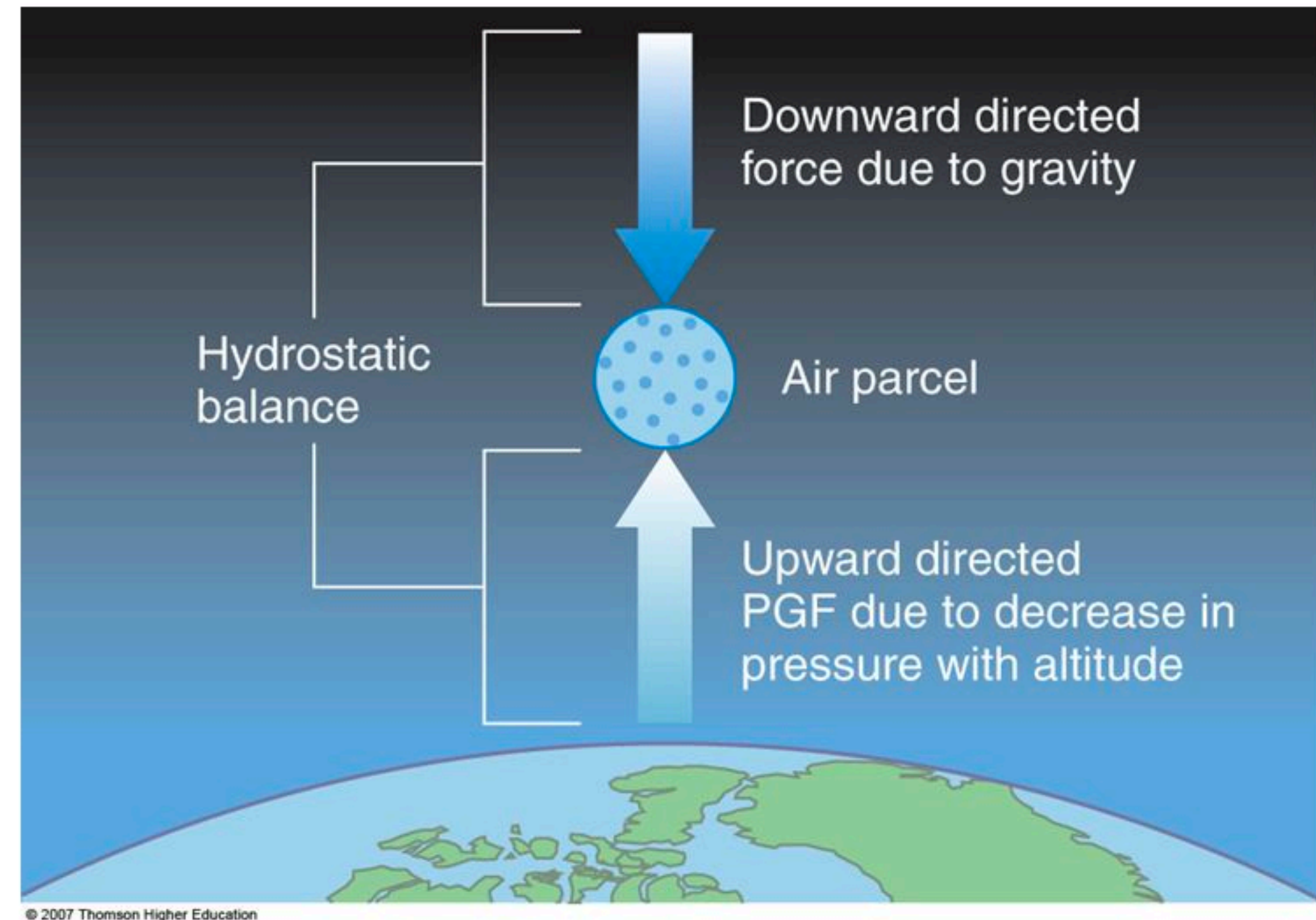
Ignoring the effects of planetary rotation and friction, the two main forces that cause vertical acceleration are gravity and the pressure gradient force.

$$\frac{Dw}{Dt} = -\alpha \frac{\partial p}{\partial z} - g$$

**Acceleration**

**Pressure  
gradient force**

**Gravity**



# Hydrostatic Balance

For quiescent atmospheric conditions, the atmosphere is maintained in place by a balance between the **downward gravitational force** and the **upward pressure gradient force**.

$$\rho g \simeq - \frac{\partial p}{\partial z}$$

Using ideal gas law we can obtain profiles for how pressure changes with height

$$p\alpha = R_d T_v$$

## Hydrostatic Equilibrium

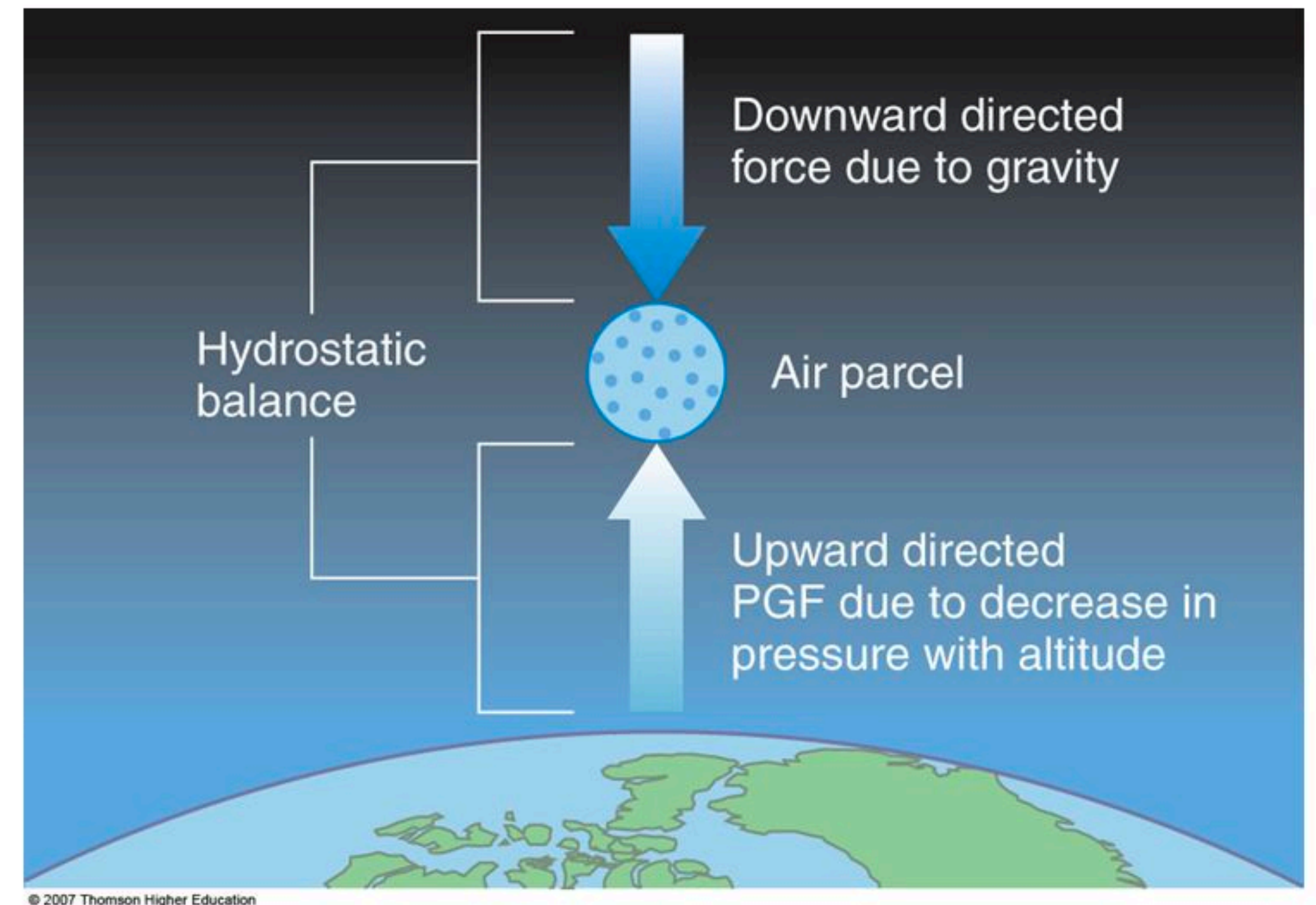


Fig. 6-13, p. 171

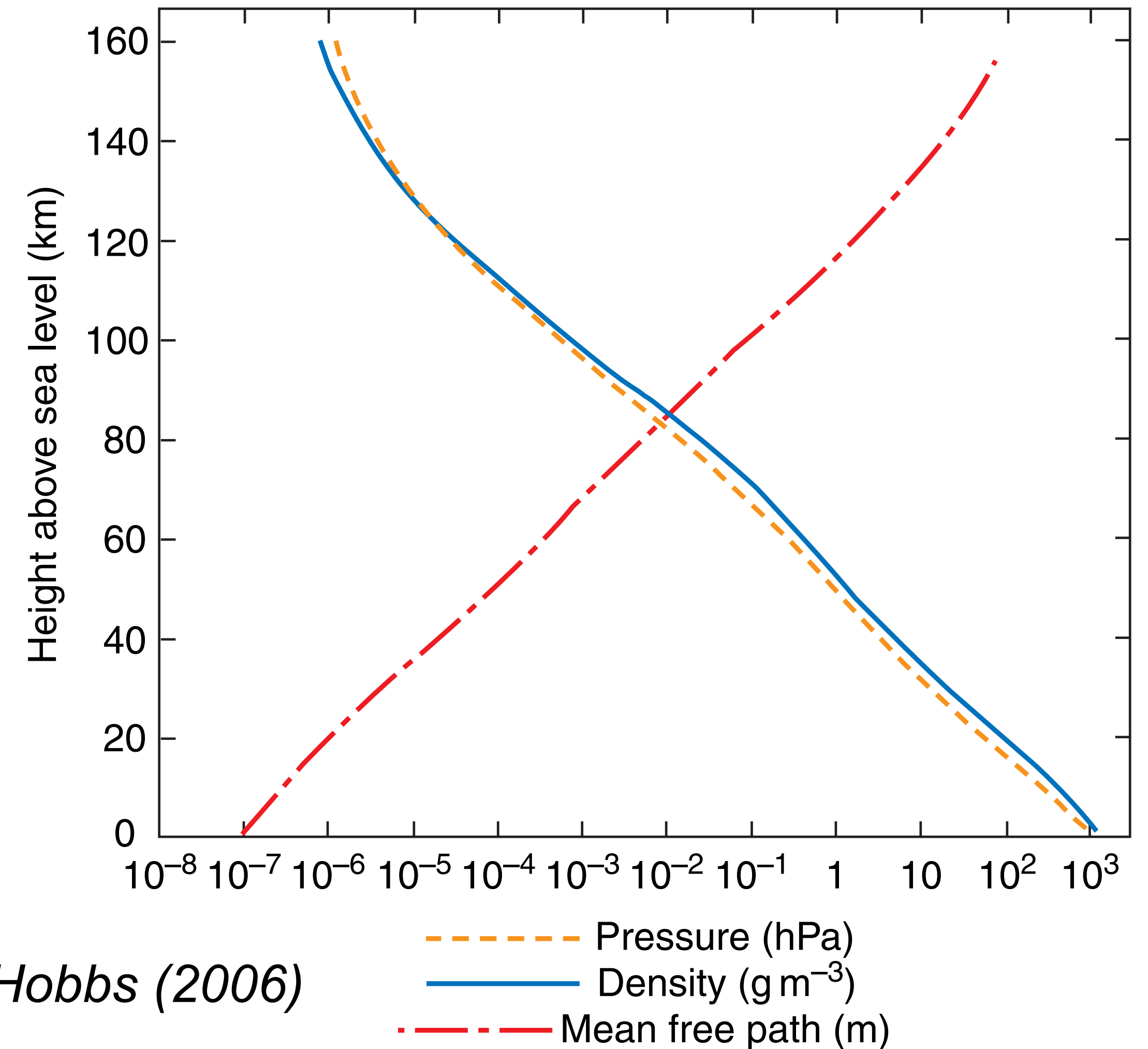


# Example

For an isothermal atmosphere, can easily solve the differential equation to obtain

$$p \simeq p_0 \exp\left(-\frac{gz}{R_d \bar{T}_v}\right)$$

$\bar{T}_v$  Mean virtual temperature of troposphere



*Wallace and Hobbs (2006)*

# Hypsometric equation

For smaller layers of the atmosphere, we can also solve the hydrostatic equation to obtain the “thickness” equation

$$Z_2 - Z_1 = \frac{R_d}{g_0} \int_{p_2}^{p_1} T_v \frac{dp}{p}$$

Which we can simplify by replacing the virtual temperature with its layer mean value to obtain

$$Z_2 - Z_1 = \bar{H} \ln \left( \frac{p_1}{p_2} \right) = \frac{R_d \bar{T}_v}{g_0} \ln \left( \frac{p_1}{p_2} \right)$$

Which is known as the hypsometric equation.



# How about the ocean?

$$\rho = \rho_0 \left[ 1 - \beta_T (T - T_0) + \beta_s (S - S_0) + \beta_p (p - p_0) \right]$$

To a good approximation we can treat seawater as incompressible

$$\rho \simeq \rho_0$$

We can solve the hydrostatic equation to obtain the following:

$$p = p_0 + \rho_0 g z$$

